## A DEEPER LOOK 4 The energy of the bonding molecular orbital of H<sub>2</sub><sup>+</sup>

The goal is to calculate the energy of the  $\sigma$  orbital of the hydrogen molecule-ion, written as the linear combination  $\psi_{+} = N_{+}(\psi_{A} + \psi_{B})$ . In general, the energy is the expectation value of the hamiltonian operator. Therefore, the energy  $E_{+}$  of  $\psi_{+}$  is

$$E_{+} = \int \psi_{+}^{*} \hat{H} \psi_{+} \mathrm{d}\tau$$

with the hamiltonian given by eqn 9B.1 and  $N_+ = 1/\{2(1+S)\}^{1/2}$ . Begin by writing a general expression for  $E_+$ , and then refine it by recognizing that  $\psi_A$  and  $\psi_B$  are H1s atomic orbitals.

**Step 1** *Write a general expression for the expectation value* Begin by using eqn 9B.1 to write

$$\hat{H}\psi_{+} = \left\{-\frac{\hbar^{2}}{2m_{\rm e}}\nabla_{1}^{2} - j_{0}\left(\frac{1}{r_{\rm A1}} + \frac{1}{r_{\rm B1}} - \frac{1}{R}\right)\right\}\psi_{+}$$
$$= \left\{-\frac{\hbar^{2}}{2m_{\rm e}}\nabla_{1}^{2} - \frac{j_{0}}{r_{\rm A1}} - \frac{j_{0}}{r_{\rm B1}}\right\}\psi_{+} + \frac{j_{0}}{R}\psi_{+}$$

where  $j_0 = e^2/4\pi\varepsilon_0$ . Because the wavefunction is real, the expectation value of the hamiltonian is

$$E_{+} = \int \psi_{+} \hat{H} \psi_{+} d\tau$$
Evaluated in step 2
$$= \int \psi_{+} \left\{ -\frac{\hbar^{2}}{2m_{e}} \nabla_{1}^{2} - \frac{j_{0}}{r_{A1}} - \frac{j_{0}}{r_{B1}} \right\} \psi_{+} d\tau + \int \psi_{+} \left( \frac{j_{0}}{R} \right) \psi_{+} d\tau$$

## **Step 2** Evaluate the first term

With  $\psi_+ = N_+(\psi_A + \psi_B)$ , and after some algebra, the first term in the expression from Step 1 expands to

$$\int \Psi_{+} \left\{ -\frac{\hbar^{2}}{2m_{e}} \nabla_{1}^{2} - \frac{j_{0}}{r_{A1}} - \frac{j_{0}}{r_{B1}} \right\} \Psi_{+} d\tau$$

$$= N_{+}^{2} \int (\Psi_{A} + \Psi_{B}) \left\{ -\frac{\hbar^{2}}{2m_{e}} \nabla_{1}^{2} - \frac{j_{0}}{r_{A1}} - \frac{j_{0}}{r_{B1}} \right\} (\Psi_{A} + \Psi_{B}) d\tau$$

$$= N_{+}^{2} \left\{ \underbrace{\int \Psi_{A} \left\{ -\frac{\hbar^{2}}{2m_{e}} \nabla_{1}^{2} - \frac{j_{0}}{r_{A1}} \right\} \Psi_{A} d\tau}_{F_{B1}} - \underbrace{\int \Psi_{A}^{2} d\tau}_{F_{B1}} d\tau$$

$$+ \underbrace{\int \psi_{\mathrm{A}} \left\{ -\frac{\hbar^2}{2m_{\mathrm{e}}} \nabla_{\mathrm{I}}^2 - \frac{j_0}{r_{\mathrm{BI}}} \right\} \psi_{\mathrm{B}} \mathrm{d}\tau}_{j_0} \underbrace{\int \frac{\psi_{\mathrm{A}} \psi_{\mathrm{B}}}{r_{\mathrm{AI}}} \mathrm{d}\tau}_{r_{\mathrm{AI}}}$$

$$+ \underbrace{\int \psi_{\rm B} \left\{ -\frac{\hbar^2}{2m_{\rm e}} \nabla_{\rm I}^2 - \frac{j_0}{r_{\rm A1}} \right\} \psi_{\rm A} d\tau}_{+ \int \psi_{\rm B} \left\{ -\frac{\hbar^2}{2m_{\rm e}} \nabla_{\rm I}^2 - \frac{j_0}{r_{\rm B1}} \right\} \psi_{\rm B} d\tau} \underbrace{\frac{k}{j_0 \int \frac{\psi_{\rm B} \psi_{\rm A}}{r_{\rm B1}} d\tau}}_{+ \int \psi_{\rm B} \left\{ -\frac{\hbar^2}{2m_{\rm e}} \nabla_{\rm I}^2 - \frac{j_0}{r_{\rm B1}} \right\} \psi_{\rm B} d\tau} - \underbrace{j_0 \int \frac{\psi_{\rm B}}{r_{\rm A1}} d\tau}_{j_0 \int \frac{\psi_{\rm B}}{r_{\rm A1}} d\tau}$$

To see how the integrals give rise to the quantities in red, consider the following:

• The terms  $-(\hbar^2/2m_e)\nabla_1^2 - j_0/r_{A1}$  and  $-(\hbar^2/2m_e)\nabla_1^2 - j_0/r_{B1}$  have the form of the hamiltonian of a hydrogen atom (Topic 8A). Because in this case  $\psi_A$  and  $\psi_B$  are H1s orbitals centred on A and B, respectively, it follows that

$$\int \psi_{\rm A} \left\{ -\frac{\hbar^2}{2m_{\rm e}} \nabla_1^2 - \frac{j_0}{r_{\rm A1}} \right\} \psi_{\rm A} \, \mathrm{d}\tau = E_{\rm 1s} \int \psi_{\rm A}^2 \mathrm{d}\tau = E_{\rm 1s}$$

$$\int \psi_{\rm B} \left\{ -\frac{\hbar^2}{2m_{\rm e}} \nabla_{\rm l}^2 - \frac{j_0}{r_{\rm Bl}} \right\} \psi_{\rm B} \, \mathrm{d}\tau = E_{\rm ls} \int \psi_{\rm B}^2 \mathrm{d}\tau = E_{\rm ls}$$

$$\int \psi_{\rm A} \underbrace{\left\{-\frac{\hbar^2}{2m_{\rm e}}\nabla_{\rm I}^2 - \frac{j_0}{r_{\rm BI}}\right\}\psi_{\rm B}}_{\xi} d\tau = E_{\rm Is} \underbrace{\int \psi_{\rm A}\psi_{\rm B} d\tau}_{\xi}$$

$$\int \psi_{\rm B} \left\{ -\frac{\hbar^2}{2m_{\rm e}} \nabla_{\rm I}^2 - \frac{j_0}{r_{\rm AI}} \right\} \psi_{\rm A} \, \mathrm{d}\tau = E_{\rm Is} \int \psi_{\rm B} \psi_{\rm A} \, \mathrm{d}\tau$$

where  $E_{1s}$  is the energy of the H1s orbital, and S is the overlap integral.

• Because the atoms are identical

$$j_0 \int \frac{\psi_A^2}{r_{B1}} d\tau = j_0 \int \frac{\psi_B^2}{r_{A1}} d\tau = j \text{ and}$$
$$j_0 \int \frac{\psi_A \psi_B}{r_{A1}} d\tau = j_0 \int \frac{\psi_B \psi_A}{r_{B1}} d\tau = k$$

Therefore, in terms of the parameters  $E_{1s}$ , S, j, and k, the first term simplifies to

$$\int \psi_{+} \left\{ -\frac{\hbar^{2}}{2m_{e}} \nabla_{1}^{2} - \frac{j_{0}}{r_{A1}} - \frac{j_{0}}{r_{B1}} \right\} \psi_{+} d\tau = 2N_{+}^{2} \{ E_{1s}(1+S) - (j+k) \}$$

## **Step 3** Evaluate the second term

The second term in the expression from Step 1 is the contribution to the energy from the repulsion between nuclei. Because  $j_0/R$  is a constant and the wavefunction is normalized, that term simplifies to

$$\int \psi_{+} \left( \frac{j_0}{R} \right) \psi_{+} \mathrm{d}\tau = \frac{j_0}{R} \int \psi_{+}^2 \mathrm{d}\tau = \frac{j_0}{R}$$

## **Step 4** *Combine the expressions from Steps 2 and 3*

Because H1s orbitals were used to arrive at the result in Step 2, write  $E_+ = E_{\sigma}$ , the energy of the  $\sigma$  molecular orbital and

$$E_{\sigma} = 2N_{+}^{2} \{E_{1s}(1+S) - (j+k)\} + \frac{\overline{j_{0}}}{R}$$

With  $N_{+} = 1/\{2(1 + S)\}^{1/2}$ , it follows that

$$E_{\sigma} = \frac{\frac{2N_{+}^{2}}{1}}{(1+S)} \{E_{1s}(1+S) - (j+k)\} + \frac{j_{0}}{R} = E_{1s} - \frac{j+k}{1+S} + \frac{j_{0}}{R}$$

as in eqn 9B.4.