

# A DEEPER LOOK 10 Establishing the relation between bulk and molecular properties

The isothermal compressibility is defined as  $\kappa_T = -(\partial V/\partial p)_T/V$ , and the bulk modulus  $K$  is simply  $1/\kappa_T$ , so it follows that  $K = -V(\partial p/\partial V)_T$ . The following calculation supposes that  $T$  is fixed at 0, so  $T$  is no longer a variable and the partial derivative becomes a complete derivative:  $K = -V(dp/dV)$  at  $T = 0$ .

## Step 1 Express the pressure in terms of the internal energy

You need to start from the fundamental thermodynamic equation  $dU = TdS - pdV$  (this is eqn 3E.1 of Topic 3E), and impose the condition  $T = 0$ ; it follows that  $p = -dU/dV$  at  $T = 0$ . This relation is now combined with the definition of  $K$  to give

$$K = -V \frac{dp}{dV} = V \frac{d^2U}{dV^2} \quad (\text{at } T = 0)$$

This expression shows that the bulk modulus (and through eqn 15D.3, the other two moduli) depends on the curvature of a plot of the internal energy against volume.

## Step 2 Relate the change in volume to the change in a lattice parameter

Now express the variation of internal energy with volume in terms of its variation with a lattice parameter,  $a$ , such as the length of the side of a unit cell:

$$\frac{dU}{dV} = \frac{dU}{da} \frac{da}{dV}$$

Then

$$\begin{aligned} \frac{d^2U}{dV^2} &= \frac{d}{dV} \left( \frac{dU}{da} \frac{da}{dV} \right) \\ &= \frac{dU}{da} \left( \frac{d^2a}{dV^2} \right) + \left( \frac{d^2U}{dVda} \right) \frac{da}{dV} \end{aligned}$$

Now use

$$\begin{aligned} \frac{d^2U}{dVda} &= \frac{d}{da} \left( \frac{dU}{dV} \right) = \frac{d}{da} \left( \frac{dU}{da} \frac{da}{dV} \right) \\ &= \left( \frac{d^2U}{da^2} \right) \left( \frac{da}{dV} \right) + \left( \frac{dU}{da} \right) \frac{d^2a}{dVda} \end{aligned}$$

and so obtain

$$\begin{aligned} \frac{d^2U}{dV^2} &= \left( \frac{dU}{da} \right) \left( \frac{d^2a}{dV^2} \right) + \left( \frac{d^2U}{da^2} \right) \left( \frac{da}{dV} \right)^2 \\ &\quad + \left( \frac{dU}{da} \right) \overbrace{\left( \frac{d^2a}{dVda} \right) \left( \frac{da}{dV} \right)}^{(d^2a/dV^2)} \\ &= 2 \left( \frac{dU}{da} \right) \left( \frac{d^2a}{dV^2} \right) + \left( \frac{d^2U}{da^2} \right) \left( \frac{da}{dV} \right)^2 \end{aligned}$$

To calculate  $K$  at the equilibrium volume of the sample,  $a$  is set to its equilibrium value,  $a_{\text{eq}}$ , and recognize that, at equilibrium,  $dU/da = 0$ . Therefore, the first (blue) term on the right is zero leaving

$$\frac{d^2U}{dV^2} = \left( \frac{d^2U}{da^2} \right) \left( \frac{da}{dV} \right)^2$$

This relation now implies that, at  $T = 0$ ,

$$K = V \left( \frac{d^2U}{dV^2} \right)_0 = V \left( \frac{d^2U}{da^2} \right)_0 \left( \frac{da}{dV} \right)_0^2 \quad (1)$$

where the subscript 0 denotes that the derivatives are evaluated at the equilibrium dimensions of the unit cell by setting  $a = a_{\text{eq}}$  after the derivative has been calculated. At this stage, write the volume  $V$  as  $V = a_{\text{eq}}^3$ ; it follows that  $da/dV = 1/3a^2$ , so

$$K = \frac{V}{(3a_{\text{eq}}^2)^2} \left( \frac{d^2U}{da^2} \right)_0 = \frac{1}{9a_{\text{eq}}} \left( \frac{d^2U}{da^2} \right)_0$$

## Step 3 Relate the internal energy to the intermolecular potential energy

If the molecules in the solid interact by a pairwise Lennard-Jones (12,6)-potential energy  $E_{\text{LJ}} = 4\epsilon\{(R_0/R)^{12} - (R_0/R)^6\}$  where  $R$  is the internuclear distance between atoms (eqn 14B.12 in Topic 14B), it follows that the internal energy of the solid is  $U = nN_{\text{A}}E_{\text{LJ}}$ , where  $n$  is the amount of substance. The size of the unit cell scales with  $R$ , so  $a = kR$ , where  $k$  is a constant that depends on the type of unit cell; similarly,  $a_{\text{eq}} = kR_{\text{eq}}$ . The minimum potential energy is at  $R_{\text{eq}} = 2^{1/6}R_0$ . With these points in mind, take the second derivative of  $U$  with respect to  $a$  and obtain

$$\begin{aligned} \frac{d^2U}{da^2} &= \frac{1}{k^2} \frac{d^2U}{dR^2} = \frac{4\epsilon nN_{\text{A}}}{k^2} \frac{d}{dR} \left\{ -12 \frac{R_0^{12}}{R^{13}} + 6 \frac{R_0^6}{R^7} \right\} \\ &= \frac{4\epsilon nN_{\text{A}}}{k^2} \left\{ 12 \times 13 \frac{R_0^{12}}{R^{14}} - 6 \times 7 \frac{R_0^6}{R^8} \right\} \end{aligned}$$

Now set  $R$  to its equilibrium value:

$$\begin{aligned} \left( \frac{d^2U}{da^2} \right)_0 &= \frac{4\epsilon nN_{\text{A}}}{k^2} \left\{ 12 \times 13 \frac{R_0^{12}}{R_{\text{eq}}^{14}} - 6 \times 7 \frac{R_0^6}{R_{\text{eq}}^8} \right\} \\ &= \frac{4\epsilon nN_{\text{A}}}{k^2} \left\{ \frac{12 \times 13}{2^{7/3} R_0^2} - \frac{6 \times 7}{2^{4/3} R_0^2} \right\} = \frac{72\epsilon nN_{\text{A}}}{2^{1/3} k^2 R_0^2} \end{aligned}$$

In terms of the equilibrium internuclear distance  $R_{\text{eq}}$ ,

$$\left(\frac{d^2U}{da^2}\right)_0 = \frac{72\epsilon n N_A}{2^{1/3} k^2 (2^{-1/6} R_{\text{eq}})^2} = \frac{72\epsilon n N_A}{k^2 R_{\text{eq}}^2}$$

It then follows from eqn 1 that

$$K = \frac{1}{9a_{\text{eq}}} \times \underbrace{\frac{72\epsilon n N_A}{k^2 R_{\text{eq}}^2}}_{a_{\text{eq}}^2} = \frac{8\epsilon n N_A}{\underbrace{a_{\text{eq}}^3}_V} = \frac{8\epsilon n N_A}{V}$$

The molar volume is  $V_m = V/n$ , so finally

$$K = \frac{8\epsilon N_A}{V_m}$$

That is, the bulk modulus is large if the interaction between molecules is strong and the molar volume is small.